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Analysis and adaptive control of a novel 3-D conservative no-equilibrium chaotic system

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First, this paper announces a seven-term novel 3-D conservative chaotic system with four quadratic nonlinearities. The conservative chaotic systems are characterized by the important property that they are volume conserving. The phase portraits of the novel conservative chaotic system are displayed and the mathematical properties are discussed. An important property of the proposed novel chaotic system is that it has no equilibrium point. Hence, it displays hidden chaotic attractors. The Lyapunov exponents of the novel conservative chaotic system are obtained as $L_1 = 0.0395$, $L_2 = 0$ and $L_3 = -0.0395$. The Kaplan-Yorke dimension of the novel conservative chaotic system is $D_{KY} = 3$. Next, an adaptive controller is designed to globally stabilize the novel conservative chaotic system with unknown parameters. Moreover, an adaptive controller is also designed to achieve global chaos synchronization of the identical conservative chaotic systems with unknown parameters. MATLAB simulations have been depicted to illustrate the phase portraits of the novel conservative chaotic system control results.

Key words: chaos, chaotic system, conservative chaotic system, adaptive control, synchronization.

1. Introduction

Chaos theory describes the qualitative study of unstable aperiodic behavior in deterministic nonlinear dynamical systems. A dynamical system is called *chaotic* if it satisfies the three properties: boundedness, infinite recurrence and sensitive dependence on initial conditions [1].

A significant development in chaos theory occurred when Lorenz discovered a 3-D chaotic system of a weather model [2]. Subsequently, Rössler discovered a 3-D chaotic system in 1976 [3], which is algebraically simpler than the Lorenz system. Indeed, Lorenz's system is a seven-term chaotic system with two quadratic nonlinearities, while Rössler's system is a seven-term chaotic system with just one quadratic nonlinearity.

Some well-known paradigms of 3-D chaotic systems are Arneodo system [4], Sprott systems [5], Chen system [6], Lü-Chen system [7], Liu system [8], Cai system [9], T-

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system [10], etc. Many new chaotic systems have been also discovered like Li system [11], Sundarapandian systems [12, 13], Vaidyanathan systems [14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26], Pehlivan system [27], Jafari system [28], Sampath system [29], Pham systems [30, 31], etc.

Chaos theory has applications in several fields of science and engineering such as lasers [32], oscillators [33], chemical reactions [34, 35], biology [36], ecology [37], neural networks [38, 39], robotics [40, 41], fuzzy logic [42, 43], electrical circuits [44, 45, 46], cryptosystems [47, 48], memristors [49, 50, 51],etc.

In the chaos literature, there is an active interest in the discovery of conservative chaotic systems [52], which have the special property that the volume of the flow is conserved. If the sum of the Lyapunov exponents of a chaotic system is zero, then the system is conservative. Classical examples of conservative chaotic systems are Nosé-Hoover system [53], Hénon-Heiles system [54], etc. A recent example of a conservative chaotic system is Vaidyanathan-Pakiriswamy system [55].

In this paper, we announce a novel 3-D conservative chaotic system which does not possess any equilibrium point. Thus, the novel chaotic system belongs to the class of chaotic systems with hidden attractors [56]. Studying systems with hidden attractors is a new research direction because of their practical and theoretical importance [57, 58].

Next, this paper derives an adaptive control law that stabilizes the novel conservative chaotic system, when the system parameters are unknown. This paper also derives an adaptive control law that achieves global chaos synchronization of the identical novel conservative systems with unknown parameters.

Synchronization of chaotic systems is a phenomenon that may occur when a chaotic oscillator drives another chaotic oscillator. Because of the butterfly effect which causes the exponential divergence of the trajectories of two identical chaotic systems started with nearly the same initial conditions, synchronizing two chaotic systems is seemingly a very challenging problem.

In most of the synchronization approaches, the master-slave or drive-response formalism is used. If a particular chaotic system is called the master or drive system and another chaotic system is called the slave or response system, then the idea of synchronization is to use the output of the master system to control the response of the slave system so that the slave system tracks the output of the master system asymptotically.

In the chaos literature, an impressive variety of techniques have been proposed to solve the problem of chaos synchronization such as active control method [59, 60, 61], adaptive control method [62, 63, 64], backstepping control method [66, 67, 68], sliding mode control method [69, 70, 71], etc.

All the main adaptive results in this paper are proved using Lyapunov stability theory [72]. MATLAB simulations are depicted to illustrate the phase portraits of the novel conservative chaos system, adaptive stabilization and synchronization results for the novel 3-D conservative chaotic system.



2. A seven-term 3-D novel conservative chaotic system

In this section, we describe a seven-term novel conservative chaotic system with four quadratic nonlinearities, which is modeled by the 3-D dynamics

$$\dot{x}_1 = ax_2 + x_1 x_3
\dot{x}_2 = -bx_1 + x_2 x_3
\dot{x}_3 = 1 - x_1^2 - x_2^2$$
(1)

where x_1, x_2, x_3 are the states and a, b are constant, positive, parameters of the system.

The system (1) exhibits a conservative chaotic attractor for the values

$$a = 0.05$$
 and $b = 1$. (2)

For numerical simulations, we take the initial conditions of the state x(t) as $x_1(0) = -1$, $x_2(0) = -1$ and $x_3(0) = 4$.

Fig. 1 shows the 3-D phase portrait of the conservative chaotic attractor of the system (1). Figs. 2–4 show the 2-D projection of the chaotic attractor of the system (1) on $(x_1,x_2), (x_2,x_3)$ and (x_1,x_3) planes, respectively.



Figure 1: Phase portrait of the conservative chaotic System



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Figure 2: 2-D projection of the conservative chaotic system on the (x_1, x_2) plane



Figure 3: 2-D projection of the conservative chaotic system on the (x_2, x_3) plane



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Figure 4: 2-D projection of the conservative chaotic system on the (x_1, x_3) plane

3. Analysis of the 3-D novel conservative chaotic system

3.1. Equilibrium points

The equilibrium points of the novel system (1) are obtained by solving the equations

$$ax_2 + x_1x_3 = 0$$
 (3a)

$$-bx_1 + x_2 x_3 = 0 (3b)$$

$$1 - x_1^2 - x_2^2 = 0. (3c)$$

From (3a) and (3b), it follows that

$$x_1 x_2 x_3 = -a x_2^2 = b x_1^2$$

which gives

$$bx_1^2 + ax_2^2 = 0. (4)$$

Since a > 0 and b > 0, the only solution of (4) is given by

$$x_1 = 0, \ x_2 = 0. \tag{5}$$

Since (3c) and (5) contradict each other, there is no equilibrium point to the novel system (1).



Rotation symmetry about the *x*₃**-axis** 3.2.

We define a new set of coordinates as

$$z_1 = -x_1 z_2 = -x_2 z_3 = x_3.$$
(6)

We find that

$$\dot{z}_1 = -ax_2 - x_1x_3 = az_2 + z_1z_3 \dot{z}_2 = bx_1 - x_2x_3 = -bz_1 + z_2z_3 \dot{z}_3 = 1 - x_1^2 - x_2^2 = 1 - z_1^2 - z_2^2.$$

$$(7)$$

This shows that the 3-D novel conservative chaotic system (1) is invariant under the change of coordinates

$$(x_1, x_2, x_3) \mapsto (-x_1, -x_2, x_3).$$
 (8)

Since the transformation (8) persists for all values of the parameters, it follows that the 3-D novel conservative chaotic system (1) has rotation symmetry about the x_3 -axis and that any non-trivial trajectory must have a twin trajectory.

3.3. Invariance

It is easy to see that the x_3 -axis is invariant under the flow of the 3-D novel conservative chaotic system (1). The invariant motion along the x_3 -axis is characterized by the scalar dynamics

$$\dot{x}_3 = 1 \tag{9}$$

which is unstable.

3.4. Lyapunov exponents and Kaplan-Yorke dimension

For the parameter values given in (2), the Lyapunov exponents of the novel chaotic system (1) are calculated as

$$L_1 = 0.0395, \quad L_2 = 0, \quad L_3 = -0.0395.$$
 (10)

Clearly, the maximal Lyapunov exponent of the novel chaotic system (1) is given by $L_1 = 0.0395$, which is positive.

Since the sum of the Lyapunov exponents in (10) is zero, the novel chaotic system (1) is conservative.

The Kaplan-Yorke dimension of a chaotic system is defined as

$$D_{KY} = j + \sum_{i=1}^{j} \frac{L_i}{|L_{j+1}|}$$





where *j* is the maximum integer such that the sum of the *j* largest Lyapunov exponents is still non-negative. D_{KY} represents an upper bound for the information dimension of the system. It is easy to deduce that for the 3-D conservative chaotic system (1), the Kaplan-Yorke dimension is given by

$$D_{KY} = 3. \tag{11}$$

4. Adaptive control of the 3-D novel conservative chaotic system with unknown parameters

In this section, we use adaptive control design to derive an adaptive feedback control law for globally stabilizing the 3-D novel conservative chaotic system with unknown parameters.

Thus, we consider the 3-D novel conservative chaotic system given by

$$\begin{cases} \dot{x}_1 = ax_2 + x_1x_3 + u_1 \\ \dot{x}_2 = -bx_1 + x_2x_3 + u_2 \\ \dot{x}_3 = 1 - x_1^2 - x_2^2 + u_3. \end{cases}$$
(12)

In (12), x_1, x_2, x_3 are the states and u_1, u_2, u_3 are adaptive controls to be determined using estimates $\hat{a}(t)$ and $\hat{b}(t)$ for the unknown parameters a and b, respectively.

We consider the adaptive control law defined by

$$\begin{cases}
 u_1 = -\hat{a}(t)x_2 - x_1x_3 - k_1x_1 \\
 u_2 = \hat{b}x_1 - x_2x_3 - k_2x_2 \\
 u_3 = -1 + x_1^2 + x_2^2 - k_3x_3
 \end{cases}$$
(13)

where k_1, k_2, k_3 are positive gain constants.

Substituting (13) into (12), we get the closed-loop plant dynamics as

$$\begin{cases} \dot{x}_1 = [a - \hat{a}(t)] x_2 - k_1 x_1 \\ \dot{x}_2 = -[b - \hat{b}(t)] x_1 - k_2 x_2 \\ \dot{x}_3 = -k_3 x_3. \end{cases}$$
(14)

The parameter estimation errors are defined as

$$\begin{cases} e_a(t) = a - \hat{a}(t) \\ e_b(t) = b - \hat{b}(t). \end{cases}$$
(15)



Using (15), we can simplify (14) as

$$\begin{cases} \dot{x}_1 = e_a x_2 - k_1 x_1 \\ \dot{x}_2 = -e_b x_1 - k_2 x_2 \\ \dot{x}_3 = -k_3 x_3. \end{cases}$$
(16)

Differentiating (15) with respect to t, we obtain

$$\begin{cases} \dot{e}_a(t) = -\dot{a}(t) \\ \dot{e}_b(t) = -\dot{b}(t). \end{cases}$$
(17)

We use adaptive control theory to find an update law for the parameter estimates. We consider the quadratic candidate Lyapunov function defined by

$$V(\mathbf{x}, e_a, e_b) = \frac{1}{2} \left(x_1^2 + x_2^2 + x_3^2 + e_a^2 + e_b^2 \right).$$
(18)

Clearly, *V* is a positive definite function on \Re^5 . Differentiating *V* along the trajectories of (16) and (17), we get

$$\dot{V} = -k_1 e_1^2 - k_2 e_2^2 - k_3 e_3^2 + e_a \left[x_1 x_2 - \dot{a} \right] + e_b \left[-x_1 x_2 - \dot{a} \right].$$
(19)

In view of (19), we take the parameter update law as follows:

$$\dot{a} = x_1 x_2$$

 $\dot{b} = -x_1 x_2.$
(20)

Next, we state and prove the main result of this section.

Theorem 7 The novel 3-D conservative chaotic system (12) with unknown system parameters is globally and exponentially stabilized for all initial conditions by the adaptive control law (13) and the parameter update law (20), where k_1, k_2, k_3 are positive gain constants.

Proof We prove this result using Lyapunov stability theory [72].

We consider the quadratic Lyapunov function defined by (18), which is a positive definite function on \Re^5 .

By substituting the parameter update law (20) into (19), we obtain the time derivative of V as

$$\dot{V} = -k_1 x_1^2 - k_2 x_2^2 - k_3 x_3^2. \tag{21}$$

From (21), it is clear that \dot{V} is a negative semi-definite function on \Re^5 .

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Thus, we can conclude that the state vector $\mathbf{x}(t)$ and the parameter estimation error are globally bounded, *i.e.*

$$\begin{bmatrix} x_1(t) & x_2(t) & x_3(t) & e_a(t) & e_b(t) \end{bmatrix}^T \in \mathcal{L}_{\infty}.$$

We define $k = \min\{k_1, k_2, k_3\}$. Then it follows from (21) that

$$\dot{V} \leqslant -k \|\boldsymbol{x}(t)\|^2. \tag{22}$$

Thus, we have

$$k\|\boldsymbol{x}(t)\|^2 \leqslant -\dot{V}.$$
(23)

Integrating the inequality (23) from 0 to t, we get

$$k \int_{0}^{t} \|\boldsymbol{x}(\tau)\|^{2} d\tau \leq V(0) - V(t).$$
(24)

From (24), it follows that $\boldsymbol{x} \in \mathcal{L}_2$.

Using (16), we can conclude that $\boldsymbol{x} \in \mathcal{L}_{\infty}$.

Using Barbalat's lemma [72], we conclude that $\mathbf{x}(t) \to 0$ exponentially as $t \to \infty$ for all initial conditions $x(0) \in \Re^3$. This completes the proof.

For the numerical simulations, the classical fourth-order Runge-Kutta method with step size $h = 10^{-8}$ is used to solve the systems (12) and (20), when the adaptive control law (13) is applied.

The parameter values of the novel conservative chaotic system are taken as in the chaotic case, *viz.* a = 0.05 and b = 1. We take the positive gain constants as $k_i = 5$ for i = 1, 2, 3.

Furthermore, as initial conditions of the novel conservative chaotic system (12), we take $x_1(0) = 7.2, x_2(0) = -5.3$ and $x_3(0) = 3.7$.

Also, as initial conditions of the parameter estimates $\hat{a}(t)$ and $\hat{b}(t)$, we take $\hat{a}(0) = 5.6$ and $\hat{b}(0) = 4.8$.

In Fig. 5, the exponential convergence of the controlled states of the 3-D conservative chaotic system (12) is depicted.

5. Adaptive synchronization of the 3-D novel conservative chaotic systems with unknown parameters

In this section, we use adaptive control method to derive an adaptive feedback control law for globally synchronizing identical 3-D novel conservative chaotic systems with unknown parameters.



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Figure 5: Time-history of the controlled states $x_1(t), x_2(t), x_3(t)$

As the master system, we consider the 3-D novel conservative chaotic system given by

$$\begin{cases} \dot{x}_1 = ax_2 + x_1 x_3 \\ \dot{x}_2 = -bx_1 + x_2 x_3 \\ \dot{x}_3 = 1 - x_1^2 - x_2^2. \end{cases}$$
(25)

In (25), x_1, x_2, x_3 are the states and a, b are unknown system parameters.

As the slave system, we consider the 3-D novel conservative chaotic system given by

$$\begin{cases} \dot{y}_1 = ay_2 + y_1y_3 + u_1 \\ \dot{y}_2 = -by_1 + y_2y_3 + u_2 \\ \dot{y}_3 = 1 - y_1^2 - y_2^2 + u_3. \end{cases}$$
(26)

In (26), y_1, y_2, y_3 are the states and u_1, u_2, u_3 are the adaptive controls to be determined using estimates $\hat{a}(t)$ and $\hat{b}(t)$ for the unknown parameters *a* and *b*, respectively.





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The synchronization error between the novel 3-D conservative chaotic systems (25) and (26) is defined by

$$\begin{cases}
e_1 = y_1 - x_1 \\
e_2 = y_2 - x_2 \\
e_3 = y_3 - x_3.
\end{cases}$$
(27)

Then the synchronization error dynamics is obtained as

$$\begin{cases} \dot{e}_1 = ae_2 + y_1y_3 - x_1x_3 + u_1 \\ \dot{e}_2 = -be_1 + y_2y_3 - x_2x_3 + u_2 \\ \dot{e}_3 = -y_1^2 + x_1^2 - y_2^2 + x_2^2 + u_3. \end{cases}$$
(28)

We consider the adaptive feedback control law

$$\begin{cases}
 u_1 = -\hat{a}(t)e_2 - y_1y_3 + x_1x_3 - k_1e_1 \\
 u_2 = \hat{b}(t)e_1 - y_2y_3 + x_2x_3 - k_2e_2 \\
 u_3 = y_1^2 - x_1^2 + y_2^2 - x_2^2 - k_3e_3
 \end{cases}$$
(29)

where k_1, k_2, k_3 are positive gain constants.

Substituting (29) into (28), we get the closed-loop error dynamics as

$$\begin{cases} \dot{e}_1 = [a - \hat{a}(t)] e_2 - k_1 e_1 \\ \dot{e}_2 = -[b - \hat{b}(t)] e_1 - k_2 e_2 \\ \dot{e}_3 = -k_3 e_3. \end{cases}$$
(30)

The parameter estimation errors are defined as

$$\begin{cases} e_a(t) = a - \hat{a}(t) \\ e_b(t) = b - \hat{b}(t). \end{cases}$$
(31)

In view of (31), we can simplify the plant dynamics (30) as

$$\begin{cases} \dot{e}_1 = e_a e_2 - k_1 e_1 \\ \dot{e}_2 = -e_b e_1 - k_2 e_2 \\ \dot{e}_3 = -k_3 e_3. \end{cases}$$
(32)

Differentiating (31) with respect to t, we obtain

$$\begin{cases} \dot{e}_a(t) = -\dot{\hat{a}}(t) \\ \dot{e}_b(t) = -\dot{\hat{b}}(t). \end{cases}$$
(33)



We use adaptive control theory to find an update law for the parameter estimates. We consider the quadratic candidate Lyapunov function defined by

$$V(\boldsymbol{e}, e_a, e_b) = \frac{1}{2} \left(e_1^2 + e_2^2 + e_3^2 + e_a^2 + e_b^2 \right).$$
(34)

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Differentiating V along the trajectories of (32) and (33), we obtain

$$\dot{V} = -k_1 e_1^2 - k_2 e_2^2 - k_3 e_3^2 + e_a \left[e_1 e_2 - \dot{a} \right] + e_b \left[-e_1 e_2 - \dot{b} \right].$$
(35)

In view of (35), we take the parameter update law as

$$\begin{cases} \dot{a}(t) = e_1 e_2 \\ \dot{b}(t) = -e_1 e_2. \end{cases}$$
(36)

Next, we state and prove the main result of this section.

Theorem 8 The novel 3-D conservative chaotic systems (25) and (26) with unknown system parameters are globally and exponentially synchronized for all initial conditions by the adaptive control law (29) and the parameter update law (36), where k_1, k_2, k_3 are positive gain constants.

Proof We prove this result by applying Lyapunov stability theory [72].

We consider the quadratic Lyapunov function defined by (34), which is clearly a positive definite function on \Re^5 .

By substituting the parameter update law (36) into (35), we obtain the time-derivative of V as

$$\dot{V} = -k_1 e_1^2 - k_2 e_2^2 - k_3 e_3^2. \tag{37}$$

From (37), it is clear that \dot{V} is a negative semi-definite function on \Re^5 .

Thus, we can conclude that the state vector $\boldsymbol{e}(t)$ and the parameter estimation error are globally bounded, *i.e.*

$$\begin{bmatrix} e_1(t) & e_2(t) & e_3(t) & e_a(t) & e_b(t) \end{bmatrix}^T \in \mathcal{L}_{\infty}$$

We define $k = \min\{k_1, k_2, k_3\}$. Then it follows from (37) that

$$\dot{V} \leqslant -k \|\boldsymbol{e}(t)\|^2. \tag{38}$$

Thus, we have

$$k\|\boldsymbol{e}(t)\|^2 \leqslant -\dot{V}.\tag{39}$$





Integrating the inequality (39) from 0 to t, we get

$$k \int_{0}^{t} \|\boldsymbol{e}(\tau)\|^{2} d\tau \leq V(0) - V(t).$$
(40)

From (40), it follows that $\boldsymbol{e} \in \mathcal{L}_2$.

Using (32), we can conclude that $\dot{\boldsymbol{e}} \in \mathcal{L}_{\infty}$.

Using Barbalat's lemma [72], we conclude that $\boldsymbol{e}(t) \to 0$ exponentially as $t \to \infty$ for all initial conditions $\boldsymbol{e}(0) \in \Re^3$. This completes the proof.

For the numerical simulations, the classical fourth-order Runge-Kutta method with step size $h = 10^{-8}$ is used to solve the systems (25) and (26) and (36), when the adaptive control law (29) is applied.

The parameter values of the novel 3-D conservative chaotic systems are taken as in the chaotic case, *viz.* a = 0.05 and b = 1. We take the positive gain constants as $k_1 = 5, k_2 = 5$ and $k_3 = 5$.

Furthermore, as initial conditions of the master system (25), we take

$$x_1(0) = 5.7, x_2(0) = 3.9, x_3(0) = -7.4.$$

As initial conditions of the slave system (26), we take

$$y_1(0) = -4.2, y_2(0) = 8.5, y_3(0) = 6.4.$$

Also, as initial conditions of the parameter estimates, we take

$$\hat{a}(0) = 8.1, \ \hat{b}(0) = 3.4.$$

Figs. 6-8 describe the complete synchronization of the 3-D novel conservative chaotic systems (25) and (26), while Fig. 9 describes the time-history of the synchronization errors e_1, e_2, e_3 .

6. Conclusions

In this research work, we detailed a seven-term novel 3-D conservative noequilibrium chaotic system with four quadratic nonlinearities. In the chaos literature, there are very few conservative no-equilibrium chaotic systems. Thus, the proposed noequilibrium conservative chaotic system is a valuable addition to the chaos literature. Next, we designed an adaptive controller to globally stabilize the novel conservative chaotic system with unknown parameters. We also designed an adaptive controller to achieve global chaos synchronization of the identical conservative chaotic systems with unknown parameters. MATLAB simulations were shown to illustrate all the main results derived in this research work.



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Figure 6: Complete synchronization of the states x_1 and y_1



Figure 7: Complete synchronization of the states x_2 and y_2







Figure 8: Complete synchronization of the states x_3 and y_3



Figure 9: Time-history of the synchronization errors e_1, e_2, e_3



References

- [1] K.T. ALLIGOOD, T.D. SAUER and J.A. YORKE: Chaos: An introduction to Dynamical Systems. New York, Springer-Verlag, 2000.
- [2] E.N. LORENZ: Deterministic nonperiodic flow. *Journal of the Atmospheric Sciences*, **20** (1963), 130–141.
- [3] O.E. RÖSSLER: An equation for continuous chaos. *Physics Letters A*, **57** (1976), 397–398.
- [4] A. ARNEODO, P. COULLET and C. TRESSER: Possible new strange attractors with spiral structure. *Communications in Mathematical Physics*, **79** (1981), 573– 579.
- [5] J.C. SPROTT: Some simple chaotic flows. *Physical Review E*, **50** (1994), 647–650.
- [6] G. CHEN and T. UETA: Yet another chaotic attractor. International Journal of Bifurcation and Chaos, 9 (1999), 1465–1466.
- [7] J. LÜ and G. CHEN: A new chaotic attractor coined. International Journal of Bifurcation and Chaos, 12 (2002), 659–661.
- [8] C.X. LIU, T. LIU, L. LIU and K. LIU: A new chaotic attractor. *Chaos, Solitons and Fractals*, 22 (2004), 1031–1038.
- [9] G. CAI and Z. TAN: Chaos synchronization of a new chaotic system via nonlinear control. *Journal of Uncertain Systems*, **1** (2007), 235-240.
- [10] G. TIGAN and D. OPRIS: Analysis of a 3D chaotic system. *Chaos, Solitons and Fractals*, 36 (2008), 1315–1319.
- [11] D. LI: A three-scroll chaotic attractor. *Physics Letters A*, **372** (2008), 387–393.
- [12] V. SUNDARAPANDIAN and I. PEHLIVAN: Analysis, control, synchronization and circuit design of a novel chaotic system. *Mathematical and Computer Modelling*, 55 (2012), 1904–1915.
- [13] V. SUNDARAPANDIAN: Analysis and anti-synchronization of a novel chaotic system via active and adaptive controllers. *Journal of Engineering Science and Technology Review*, **6** (2013), 45–52.
- [14] S. VAIDYANATHAN: A new six-term 3-D chaotic system with an exponential nonlinearity. *Far East Journal of Mathematical Sciences*, **79** (2013), 135–143.

- [15] S. VAIDYANATHAN: Analysis and adaptive synchronization of two novel chaotic systems with hyperbolic sinusoidal and cosinusoidal nonlinearity and unknown parameters. *Journal of Engineering Science and Technology Review*, 6 (2013), 53–65.
- [16] S. VAIDYANATHAN: A new eight-term 3-D polynomial chaotic system with three quadratic nonlinearities. *Far East Journal of Mathematical Sciences*, **84** (2014), 219–226.
- [17] S. VAIDYANATHAN: Analysis, control and synchronisation of a six-term novel chaotic system with three quadratic nonlinearities. *International Journal of Modelling, Identification and Control,* **22** (2014), 41–53.
- [18] S. VAIDYANATHAN and K. MADHAVAN: Analysis, adaptive control and synchronization of a seven-term novel 3-D chaotic system. *International Journal of Control Theory and Applications*, 6 (2013), 121–137.
- [19] S. VAIDYANATHAN: Analysis and adaptive synchronization of eight-term 3-D polynomial chaotic systems with three quadratic nonlinearities. *European Physical Journal: Special Topics*, **223** (2014), 1519–1529.
- [20] S. VAIDYANATHAN, CH. VOLOS, V.T. PHAM, K. MADHAVAN and B.A. ID-OWU: Adaptive backstepping control, synchronization and circuit simulation of a 3-D novel jerk chaotic system with two hyperbolic sinusoidal nonlinearities. *Archives of Control Sciences*, 24 (2014), 257–285.
- [21] S. VAIDYANATHAN: Generalised projective synchronisation of novel 3-D chaotic systems with an exponential non-linearity via active and adaptive control. *International Journal of Modelling, Identification and Control*, **22** (2014), 207–217.
- [22] S. VAIDYANATHAN: Analysis, properties and control of an eight-term 3-D chaotic system with an exponential nonlinearity. *International Journal of Modelling, Identification and Control,* **23** (2015), 164–172.
- [23] S. VAIDYANATHAN: A 3-D novel highly chaotic system with four quadratic nonlinearities, its adaptive control and anti-synchronization with unknown parameters. *Journal of Engineering Science and Technology Review*, **8** (2015), 106–115.
- [24] S. VAIDYANATHAN, K. RAJAGOPAL, C.K. VOLOS, I.M. KYPRIANIDIS and I.N. STOUBOULOS: Analysis, adaptive control and synchronization of a seventerm novel 3-D chaotic system with three quadratic nonlinearities and its digital implementation in LabVIEW. *Journal of Engineering Science and Technology Review*, 8 (2015), 130–141.



- [25] S. VAIDYANATHAN, C.K. VOLOS, I.M. KYPRIANIDIS, I.N. STOUBOULOS and V.-T. PHAM: Analysis, adaptive control and anti-synchronization of a sixterm novel jerk chaotic system with two exponential nonlinearities and its circuit simulation. *Journal of Engineering Science and Technology Review*, 8 (2015), 24– 36.
- [26] S. VAIDYANATHAN, C.K. VOLOS and V.-T. PHAM: Analysis, adaptive control and adaptive synchronization of a nine-term novel 3-D chaotic system with four quadratic nonlinearities and its circuit simulation. *Journal of Engineering Science* and Technology Review, 8 (2015), 174–184.
- [27] I. PEHLIVAN, I.M. MOROZ and S. VAIDYANATHAN: Analysis, synchronization and circuit design of a novel butterfly attractor. *Journal of Sound and Vibration*, 333 (2014), 5077–5096.
- [28] S. JAFARI and J.C. SPROTT: Simple chaotic flows with a line equilibrium. *Chaos, Solitons and Fractals*, **57** (2013), 79–84.
- [29] S. SAMPATH, S. VAIDYANATHAN, C.K. VOLOS and V.-T. PHAM: An eightterm novel four-scroll chaotic system with cubic nonlinearity and its circuit simulation. *Journal of Engineering Science and Technology Review*, 8 (2015), 1–6.
- [30] V.T. PHAM, C. VOLOS, S. JAFARI, Z. WEI and X. WANG: Constructing a novel no-equilibrium chaotic system. *International Journal of Bifurcation and Chaos*, 24 (2014), 1450073.
- [31] V.T. PHAM, S. VAIDYANATHAN, C.K. VOLOS and S. JAFARI: Hidden attractors in a chaotic system with an exponential nonlinear term. *European Physical Journal: Special Topics*, 224 (2015), 1507–1517.
- [32] S. BEHNIA, S. AFRANG, A. AKHSHANI and Kh. MABHOUTI: A novel method for controlling chaos in external cavity semiconductor laser. *Optik - International Journal for Light and Electron Optics*, **124** (2013), 757–764.
- [33] J.M. TUWANKOTTA: Chaos in a coupled oscillators system with widely spaced frequencies and energy-preserving non-linearity. *International Journal of Non-Linear Mechanics*, **41** (2006), 180–191.
- [34] S. VAIDYANATHAN: Adaptive synchronization of chemical chaotic reactors. *International Journal of ChemTech Research*, **8** (2015), 612–621.
- [35] S. VAIDYANATHAN: Adaptive control of a chemical chaotic reactor. *International Journal of PharmTech Research*, **8** (2015), 377–382.
- [36] S. VAIDYANATHAN: Adaptive backstepping control of enzymes-substrates system with ferroelectric behaviour in brain waves. *International Journal of PharmTech Research*, **8** (2015), 256–261.

- [37] B. SAHOO and S. PORIA: The chaos and control of a food chain model supplying additional food to top-predator. *Chaos, Solitons & Fractals,* **58** (2014), 52–64.
- [38] W.Z. HUANG and Y. HUANG: Chaos of a new class of Hopfield neural networks. *Applied Mathematics and Computation*, **206** (2008), 1–11.
- [39] Y. SUN, V. BABOVIC and E.S. CHAN: Multi-step-ahead model error prediction using time-delay neural networks combined with chaos theory. *Journal of Hydrology*, **395** (2010), 109–116.
- [40] M. ISLAM and K. MURASE: Chaotic dynamis of a behavior-based miniature mobile robot: effects of environment and control structure. *Neural Networks*, 18 (2005), 123–144.
- [41] CH.K. VOLOS, I.M. KYPRIANIDIS and I.N. STOUBOULOS: Experimental investigation on coverage performance of a chaotic autonomous mobile robot. *Robotics and Autonomous Systems*, **61** (2013), 1314–1322.
- [42] H.T. YAU and C.S. SHIEH: Chaos synchronization using fuzzy logic controller. *Nonlinear Analysis: Real World Applications*, **9** (2008), 1800–1810.
- [43] N.S. PAI, H.T. YAU and C.L. KUO: Fuzzy logic combining controller design for chaos control of a rod-type plasma torch system. *Expert Systems with Applications*, **37** (2010), 8278–8283.
- [44] A.E. MATOUK: Chaos, feedback control and synchronization of a fractionalorder modified autonomous Van der PolŰDuffing circuit. *Communications in Nonlinear Science and Numerical Simulation*, 16 (2011), 975–986.
- [45] CH.K. VOLOS, I.M. KYPRIANIDIS, I.N. STOUBOULOS and A.N. ANAG-NOSTOPOULOS: Experimental study of the dynamic behavior of a double scroll circuit. *Journal of Applied Functional Analysis*, **4** (2009), 703–711.
- [46] CH.K. VOLOS, V.-T. PHAM, S. VAIDYANATHAN, I.M. KYPRIANIDIS and I.N. STOUBOULOS: Synchronization phenomena in coupled Colpitts circuits. *Journal of Engineering Science and Technology Review*, 8 (2015), 142–151.
- [47] CH.K. VOLOS, I.M. KYRPIANIDIS and I.N. STOUBOULOS: Image encryption process based on chaotic synchronization phenomena. *Signal Processing*, 93 (2013), 1328–1340.
- [48] CH.K. VOLOS, I.M. KYRPIANIDIS and I.N. STOUBOULOS: Text encryption scheme realized with a chaotic pseudo-random bit generator. *Journal of Engineering Science and Technology Review*, **6** (2013), 9–14.
- [49] V.-T. PHAM, C. VOLOS, S. JAFARI, X. WANG and S. VAIDYANATHAN: Hidden hyperchaotic attractor in a novel simple memristive neural network. *Optoelectronics and Advanced Materials, Rapid Communications*, 8 (2014) 1157–1163.



- [50] V.-T. PHAM, CH.K. VOLOS, S. VAIDYANATHAN, T.P. LE and V.Y. VU: A memristor-based hyperchaotic system with hidden attractors: Dynamics, synchronization and circuital emulating. *Journal of Engineering Science and Technology Review*, **8** (2015), 205–214.
- [51] CH.K. VOLOS, I.M. KYPRIANIDIS, I.N. STOUBOULOS, E. TLELO-CUAUTLE and S. VAIDYANATHAN: Memristor: A new concept in synchronization of coupled neuromorphic circuits. *Journal of Engineering Science and Technology Review*, 8 (2015), 157–173.
- [52] J.C. SPROTT: Elegant Chaos. Singapore, World Scientific, 2010.
- [53] W.G. HOOVER: Remark on 'Some simple chaotic flows', *Phys. Rev. E*, **51** (1995), 759–760.
- [54] M. HÉNON and C. HEILES: The applicability of the third integral of motion: Some numerical experiments. *Astrophysical Journal*, **69** (1964), 73–79.
- [55] S. VAIDYANATHAN and S. PAKIRISWAMY: A 3-D novel conservative chaotic system and its generalized projective synchronization via adaptive control. *Jour*nal of Engineering Science and Technology Review, 8 (2015), 52–60.
- [56] G. LEONOV, N. KUZNETSOV, S. SELDEDZHI and V. VAGAITSEV: Hidden oscillations in dynamical systems. *Trans. Syst. Contr.*, **6** (2011), 54–67.
- [57] Z. WEI: Dynamical behaviors of a chaotic system with no equilibria. *Phys. Lett. A*, **376** (2011), 102–108.
- [58] S. JAFARI, J.C. SPROTT and S. GOLPAYEGANI: Elementary quadratic chaotic flows with no equilibria. *Phys. Lett. A*, **377** (2013), 699–702.
- [59] B.A. IDOWU, U.E. VINCENT and A.N. NJAH: Synchronization of chaos in nonidentical parametrically excited systems. *Chaos, Solitons and Fractals*, **39** (2009), 2322–2331.
- [60] S. VAIDYANATHAN and K. RAJAGOPAL: Hybrid synchronization of hyperchaotic Wang-Chen and hyperchaotic lorenz systems by active non-linear control. *International Journal of Signal System Control and Engineering Application*, 4 (2011), 55–61.
- [61] S. VAIDYANATHAN, A.T. AZAR, K. RAJAGOPAL and P. ALEXANDER: Design and SPICE implementation of a 12-term novel hyperchaotic system and its synchronisation via active control. *International Journal of Modelling, Identification and Control*, 23 (2015), 267–277.
- [62] V. SUNDARAPANDIAN and R. KARTHIKEYAN: Anti-synchronization of Lü and Pan chaotic systems by adaptive nonlinear control. *European Journal of Scientific Research*, 64 (2011), 94–106.





- [63] V. SUNDARAPANDIAN and R. KARTHIKEYAN: Adaptive antisynchronization of Uncertain Tigan and Li Systems. *Journal of Engineering and Applied Sciences*, **7** (2012), 45–52.
- [64] S. VAIDYANATHAN: Hyperchaos, qualitative analysis, control and synchronisation of a ten-term 4-D hyperchaotic system with an exponential nonlinearity and three quadratic nonlinearities. *International Journal of Modelling, Identification and Control,* **23** (2015), 380–392.
- [65] S. VAIDYANATHAN, V.-T. PHAM and C.K. VOLOS: A 5-D hyperchaotic Rikitake dynamo system with hidden attractors. *European Physical Journal: Special Topics*, 224 (2015), 1575–1592.
- [66] S. RASAPPAN and S. VAIDYANATHAN: Synchronization of hyperchaotic Liu system via backstepping control with recursive feedback. *Communications in Computer and Information Science*, **305** (2012), 212–221.
- [67] S. VAIDYANATHAN and S. RASAPPAN: Global chaos synchronization of *n*-scroll Chua circuit and Lur'e system using backstepping control design with recursive feedback. *Arabian Journal for Science and Engineering*, **39** (2014), 3351–3364.
- [68] S. VAIDYANATHAN, C.K. VOLOS, K. RAJAGOPAL, I.M. KYPRIANIDIS and I.N. STOUBOULOS: Adaptive backstepping controller design for the antisynchronization of identical WINDMI chaotic systems with unknown parameters and its SPICE implementation. *Journal of Engineering Science and Technology Review*, 8 (2015), 74–82.
- [69] S. VAIDYANATHAN and S. SAMPATH: Global chaos synchronization of hyperchaotic Lorenz systems by sliding mode control. *Communications in Computer and Information Science*, **205** (2011), 156–164.
- [70] S. VAIDYANATHAN: Global chaos synchronisation of identical Li-Wu chaotic systems via sliding mode control. *International Journal of Modelling, Identification and Control*, 22 (2014), 170–177.
- [71] S. VAIDYANATHAN, S. SAMPATH and A.T. AZAR: Global chaos synchronisation of identical chaotic systems via novel sliding mode control method and its application to Zhu system. *International Journal of Modelling, Identification and Control*, **23** (2015), 92–100.
- [72] H.K. KHALIL: Nonlinear Systems. New York, Prentice Hall, 2002.